

Quantum Computing

- Week 13 (July 16-17, 2025)
- Topics:
 - A summary of quantum algorithms
 - Pure states and mixed states
 - Density operator and trace
 - Partial trace and partial measurement
 - Reduced density operator

Quantum Algorithms

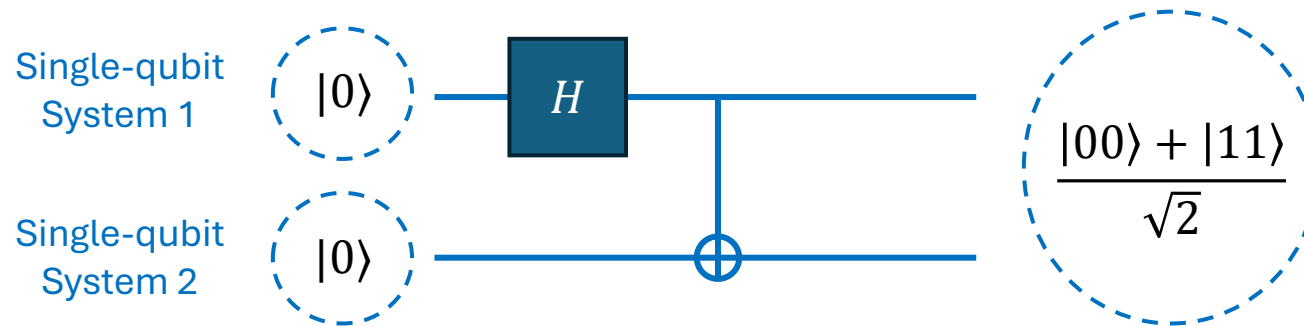
- Quantum algorithms that we have studied so far:
 - Deutsch-Jozsa algorithm
 - Superdense coding
 - Quantum teleportation
 - Quantum Fourier transformation and order finding
 - Grover search algorithm

Quantum Algorithms

Algorithm	Addressing problem	Classical “Complexity”	Quantum “Complexity”	Improvement
Deutsch-Jozsa	Balance functions	$O(2^n)$	$O(1)$	Exponentially
Superdense coding	Transmit classical info	-	1 qubit = 2-bit info (via 1 entangled pair)	-
Quantum teleportation	Teleport quantum states	2-bit info = 1 qubit (via 1 entangled pair)	-	-
Order/Period Finding (QFT)	Factoring, Discrete log	$2^{o(n)} \sim O(2^{n/2})$	$O(n^2)$ or $O(n \log n)$	Exponentially
Grover	Unstructured search	$O(2^n)$	$O(2^{n/2})$	Quadratically

Mixed States, Recaps

- **Pure state:** Can be described by a state vector
- **Mixed state:** Cannot ...



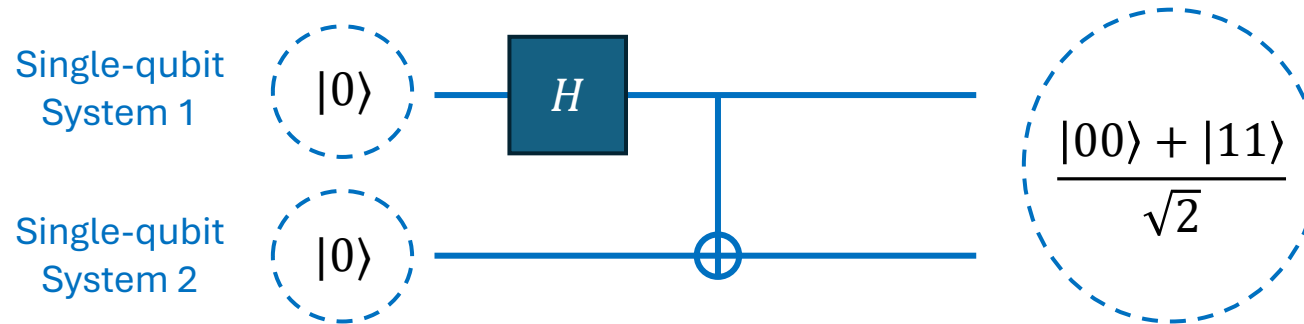
Small Exercise: (pure or mixed)

1. The initial state of system 1 is ____.
2. The states of systems 1 and 2 (after H and CNOT) are both ____.
3. The state of the total system (after H and CNOT) is ____.

Mixed States, Recaps

- **Pure state:** Can be described by a state vector
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How can we describe mixed states?



Small Exercise: (pure or mixed)

1. The initial state of system 1 is pure.
2. The states of systems 1 and 2 (after H and CNOT) are both mixed.
3. The state of the total system (after H and CNOT) is pure.

Density Operator

- Let $\{|\psi_i\rangle\}_i$ be a set of pure states, where i is an index
- Suppose that a quantum system is in $|\psi_i\rangle$ with probability p_i (s.t. $\sum_i p_i = 1$)
- Then we write the **density operator** ρ of the system as

$$\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$$

- Examples:
 - A single-qubit system with state $|0\rangle$
 - A single-qubit system with state $\frac{|0\rangle+|1\rangle}{\sqrt{2}}$
 - A single-qubit system that is in state $|0\rangle$ with probability $\frac{1}{2}$ and in state $|1\rangle$ with probability $\frac{1}{2}$

Density Operator

- $\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$ (where $\sum_i p_i = 1$)

What's the difference between the states of the two systems?

System 1: Sample two bits $b_1 b_2$ uniformly at random, and set its state as $|b_1 b_2\rangle$

System 2: $\frac{|00\rangle + |01\rangle + |10\rangle + |11\rangle}{2}$

Density Operator

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System 2: $\frac{|00\rangle + |01\rangle + |10\rangle + |11\rangle}{2}$

$$\rho_1 = \frac{1}{4} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

$$\rho_2 = \frac{1}{4} \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \end{bmatrix}$$

Density Operator

- $\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$ (where $\sum_i p_i = 1$)
- Density operator **provides another way to formulate postulates** of quantum computing:

Postulate	State vector	Density operator
Describing the state of a system	$ \psi\rangle = \sum_i \alpha_i \psi_i\rangle$	$\rho = \sum_i p_i \psi_i\rangle\langle\psi_i $
Unitary transformation	$ \psi\rangle \mapsto U \psi\rangle$	$\rho \mapsto U\rho U^\dagger$
Quantum measurement $\{M_m\}_m$	$p(m) = \langle\psi M_m^\dagger M_m \psi\rangle,$ $ \psi_m\rangle = \frac{M_m}{\sqrt{p(m)}} \psi\rangle$	$p(m) = \text{tr}(M_m^\dagger M_m \rho),$ $\rho_m = \frac{M_m \rho M_m^\dagger}{p(m)}$
Composite system	$ \psi_1\rangle \otimes \dots \otimes \psi_n\rangle$	$ \rho_1\rangle \otimes \dots \otimes \rho_n\rangle$

Trace

- The **sum of diagonal elements**:

$$\text{tr}(M) = \sum_i M_{ii}$$

- Properties of trace:
 - **Linearity:** $\text{tr}(A + B) = \text{tr}(A) + \text{tr}(B)$, $\text{tr}(z \cdot A) = z \cdot \text{tr}(A)$
 - **Cyclicity:** $\text{tr}(AB) = \text{tr}(BA)$ (similarly, $\text{tr}(ABC) = \text{tr}(CAB) = \text{tr}(BCA) \dots$)
- Several facts implied by **Cyclicity**:
 - $\text{tr}(U\rho U^\dagger) = \text{tr}(\rho)$
 - $\text{tr}(M|\psi\rangle\langle\psi|) = \text{tr}(\langle\psi|M|\psi\rangle) = \langle\psi|M|\psi\rangle$

Trace

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- Several facts implied by **Cyclicity**:
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 - $\text{tr}(\mathbf{M}|\psi\rangle\langle\psi|) = \text{tr}(\langle\psi|\mathbf{M}|\psi\rangle) = \langle\psi|\mathbf{M}|\psi\rangle$, (and thus $\text{tr}(|\psi\rangle\langle\psi|) = \langle\psi|\psi\rangle = 1$)
- Some examples:
 - Unitary transformation formulated by state vectors v.s. by density operators
 - Measurement formulated by state vectors v.s. by density operators

Trace

- **Criterion** to decide if a state is mixed or pure: Let ρ be the density operator of a quantum system.
 - $\text{tr}(\rho^2) < 1$: Mixed state
 - $\text{tr}(\rho^2) = 1$: Pure state

Partial Trace and Reduced Density Operator

- Let:
 - A and B be two quantum systems
 - Let $|a_1\rangle$ and $|a_2\rangle$ be any two state vectors defined over (the state space of) A
 - Let $|b_1\rangle$ and $|b_2\rangle$ be any two state vectors defined over B

- **Partial Trace:**

$$\text{tr}_B(|a_1\rangle\langle a_2| \otimes |b_1\rangle\langle b_2|) := |a_1\rangle\langle a_2| \cdot \text{tr}(|b_1\rangle\langle b_2|)$$

- Remark:
 1. tr maps an operator onto a complex number, but tr_B maps an operator onto an operator.
 2. The operator from tr_B is an operator defined over A

Partial Trace and Reduced Density Operator

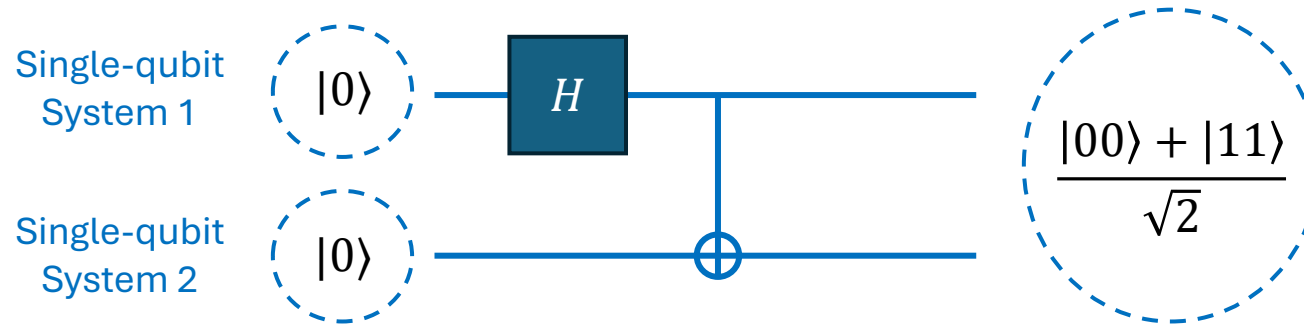
- Let:
 - A and B be two quantum systems
 - Let ρ_{AB} be the state of the state space of the composite system $A \otimes B$
- **Reduced Density Operator of A :**

$$\rho_A := \text{tr}_B(\rho_{AB})$$

- Examples (of calculating ρ_A) :
 1. $\rho_{AB} = \rho \otimes \sigma$, where ρ is defined over A and σ is defined over B
 2. $\rho_{AB} = |00\rangle\langle 00| (= |0\rangle\langle 0| \otimes |0\rangle\langle 0|)$
 3. $\rho_{AB} = \left(\left(\frac{|0\rangle + |1\rangle}{\sqrt{2}} \right) \otimes \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right) \right) \left(\left(\frac{\langle 0| + \langle 1|}{\sqrt{2}} \right) \otimes \left(\frac{\langle 0| - \langle 1|}{\sqrt{2}} \right) \right)$

Partial Trace and Reduced Density Operator

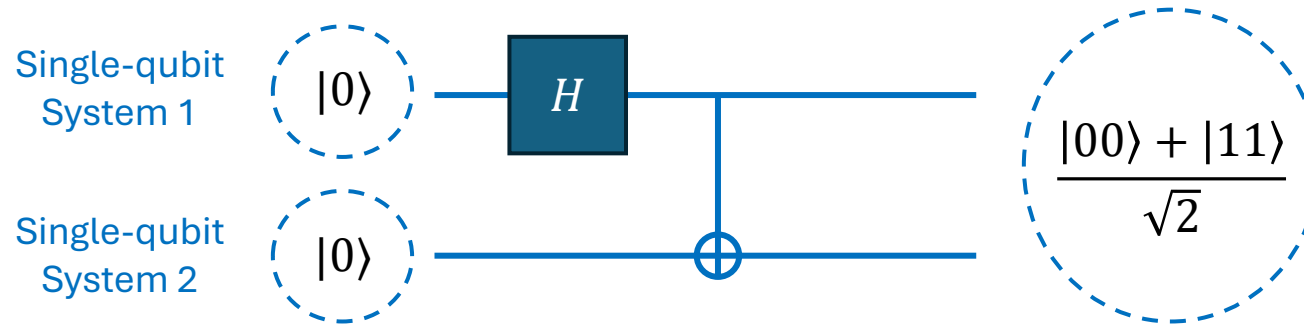
- The final state is $\rho_{12} = \left(\frac{|00\rangle + |11\rangle}{\sqrt{2}} \right) \otimes \left(\frac{|00\rangle + |11\rangle}{\sqrt{2}} \right)$



- $\rho_1 = \text{tr}_2(\rho_{12}) = ?$

Partial Trace and Reduced Density Operator

- The final state is $\rho_{12} = \left(\frac{|00\rangle + |11\rangle}{\sqrt{2}} \right) \otimes \left(\frac{|00\rangle + |11\rangle}{\sqrt{2}} \right)$



- $\rho_1 = \text{tr}_2(\rho_{12}) = \frac{1}{2} |0\rangle\langle 0| + \frac{1}{2} |1\rangle\langle 1|$

Pure state and Mixed state

- $\rho = |\psi\rangle\langle\psi|$, where $|\psi\rangle = \frac{|0\rangle+|1\rangle}{\sqrt{2}}$

- $\rho' = \frac{1}{2}|0\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1|$

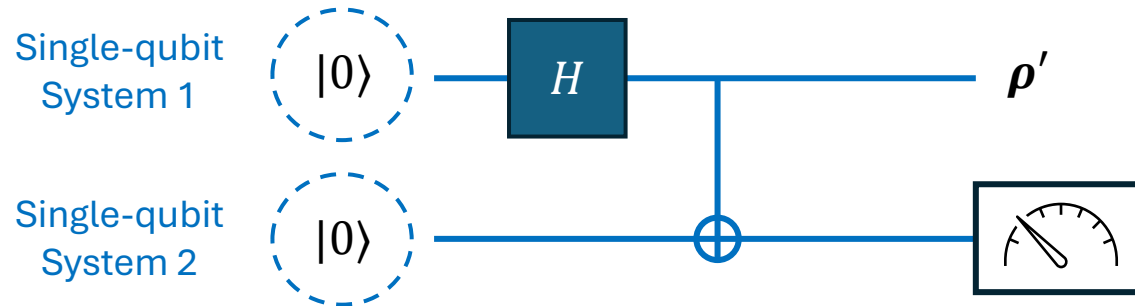
Pure state and Mixed state

- $\rho = |\psi\rangle\langle\psi|$, where $|\psi\rangle = \frac{|0\rangle+|1\rangle}{\sqrt{2}}$
 - $\text{tr}(\rho) = 1 \Rightarrow$ A pure state.
 - We “somehow” know the state of the system with **certainty** (i.e., the state is $|\psi\rangle$ with probability 1)
 - More generally, $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$, we still know the state of the system with **certainty** (though we do not know α and β , the system is still **in a definite pure state...**)
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Pure state and Mixed state

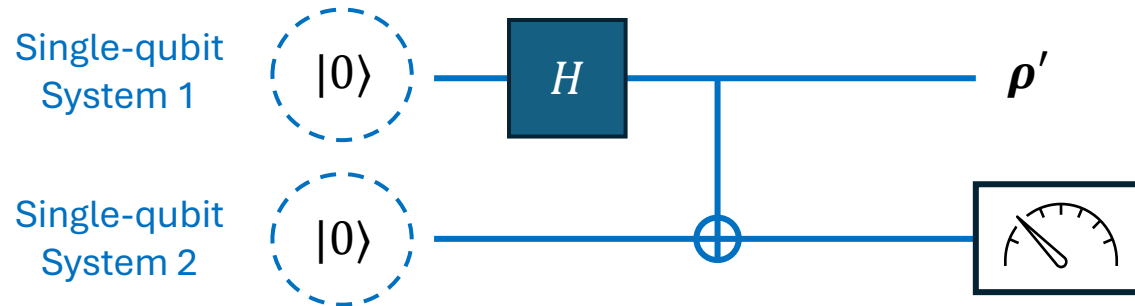
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- $\rho' = \frac{1}{2}|0\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1|$
 - $\text{tr}(\rho) = \frac{1}{2} < 1 \Rightarrow$ A mixed state
 - We are **uncertain** about the exact state the system is in...

Pure state and Mixed state



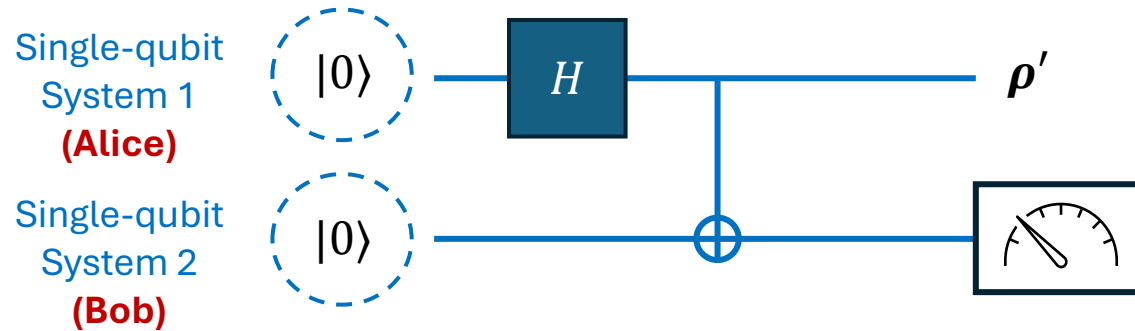
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 - We are **uncertain** about the exact state the system is in...
 - The measurement (on system 2) outcome is b with probability $\frac{1}{2}$, corresponding to the state and the probability distribution of system 1...

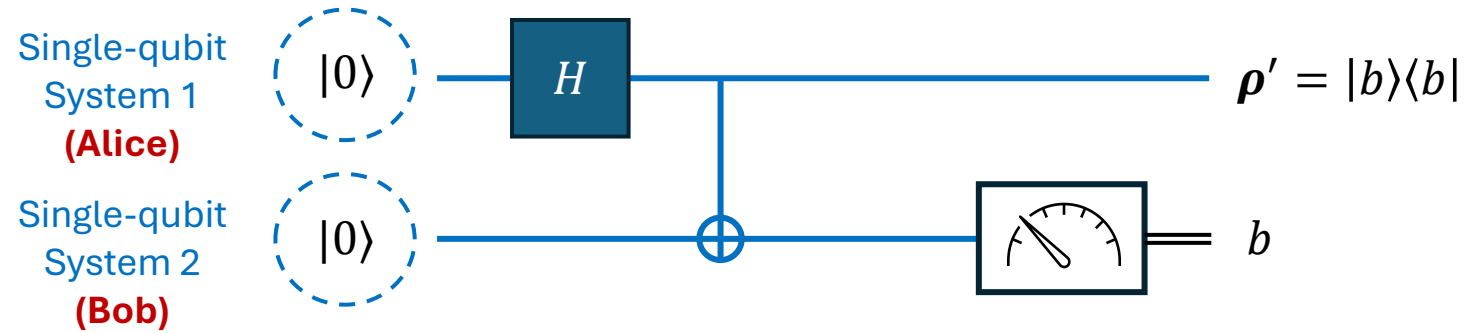
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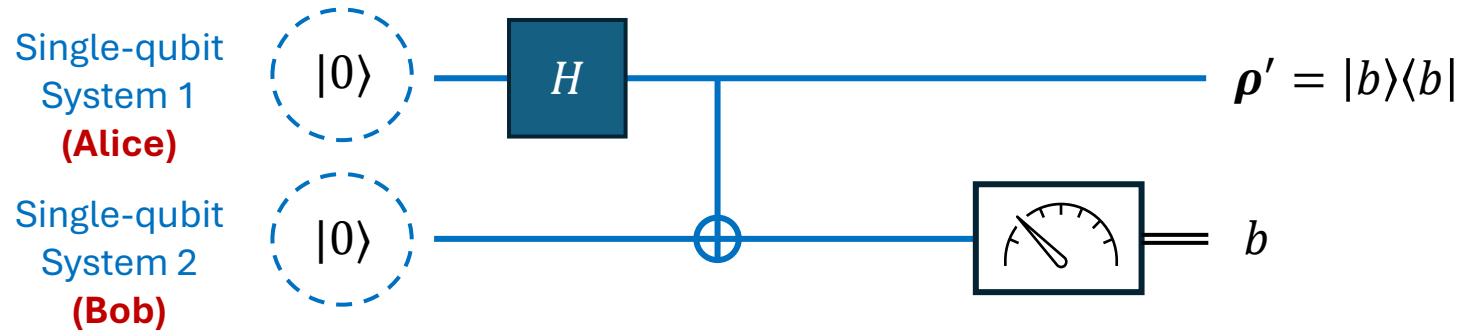
What if Bob tells the measurement outcome to Alice?

Pure state and Mixed state



- $\rho' = |b\rangle\langle b|$

Pure state and Mixed state

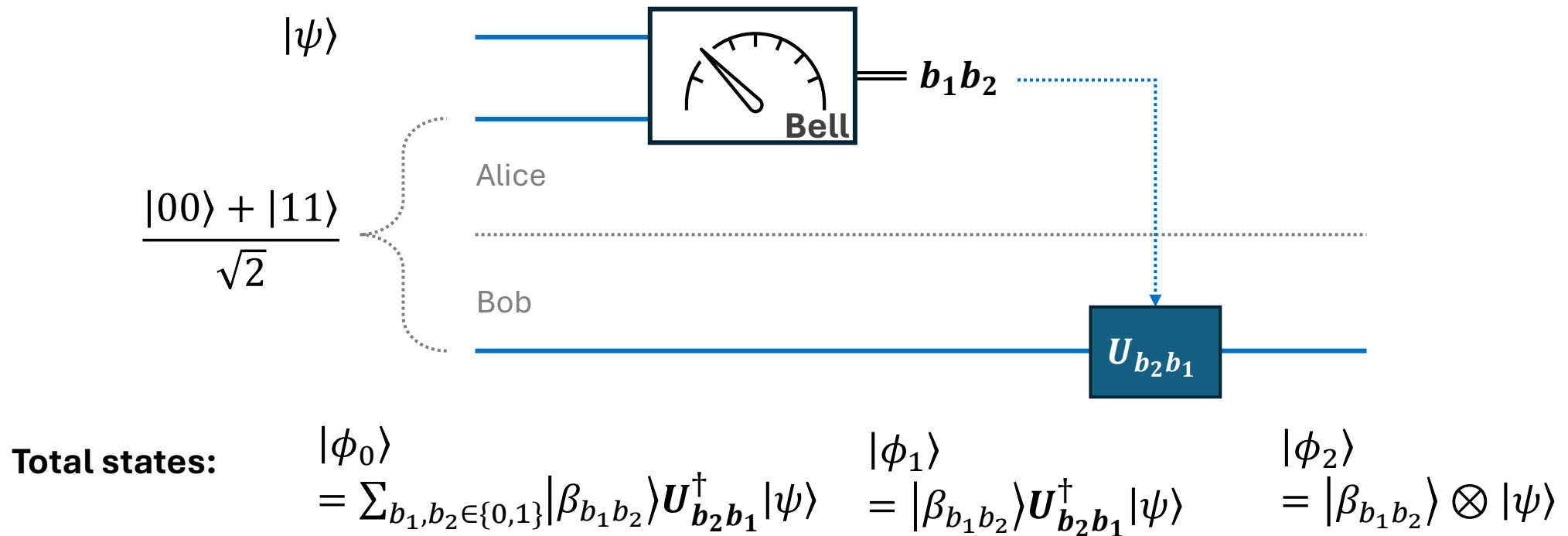


- $\rho' = |b\rangle\langle b|$

- If Bob sends b to Alice, then $\rho' = |b\rangle\langle b|$.
- If not, $\rho' = \frac{1}{2}|0\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1|$ in Alice's view.
(No information about b at all!)

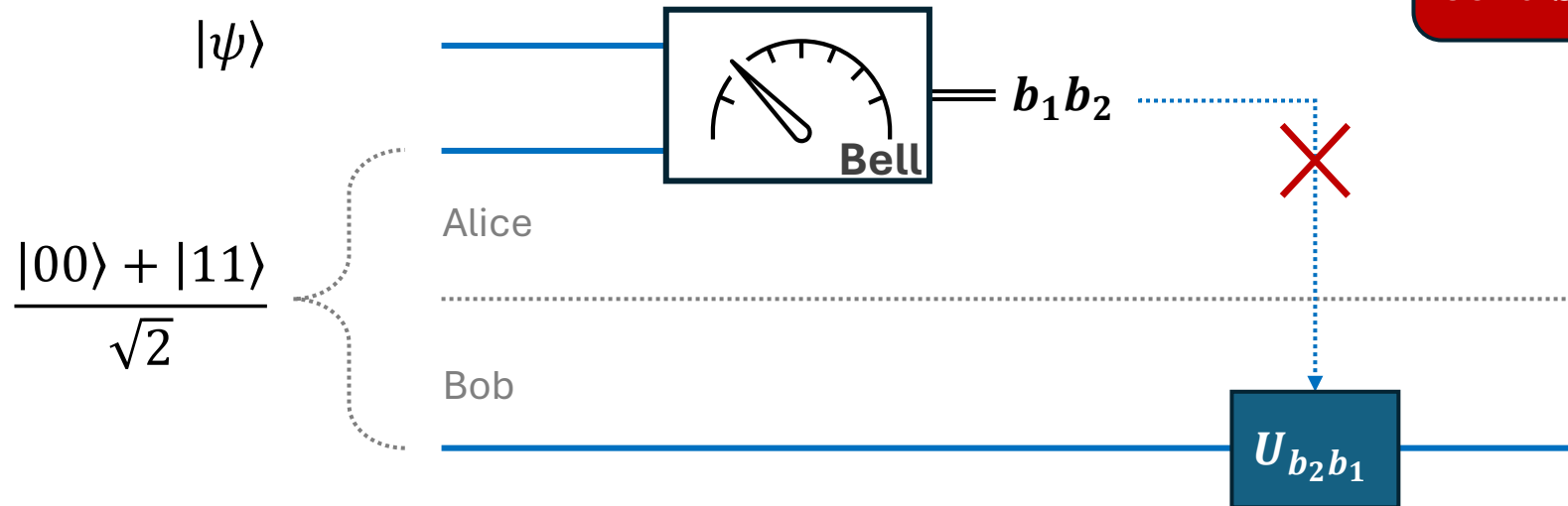
Another look at Quantum Teleportation

- Quantum teleportation: Transmit a single-qubit state $|\psi\rangle$.



Another look at Quantum Teleportation

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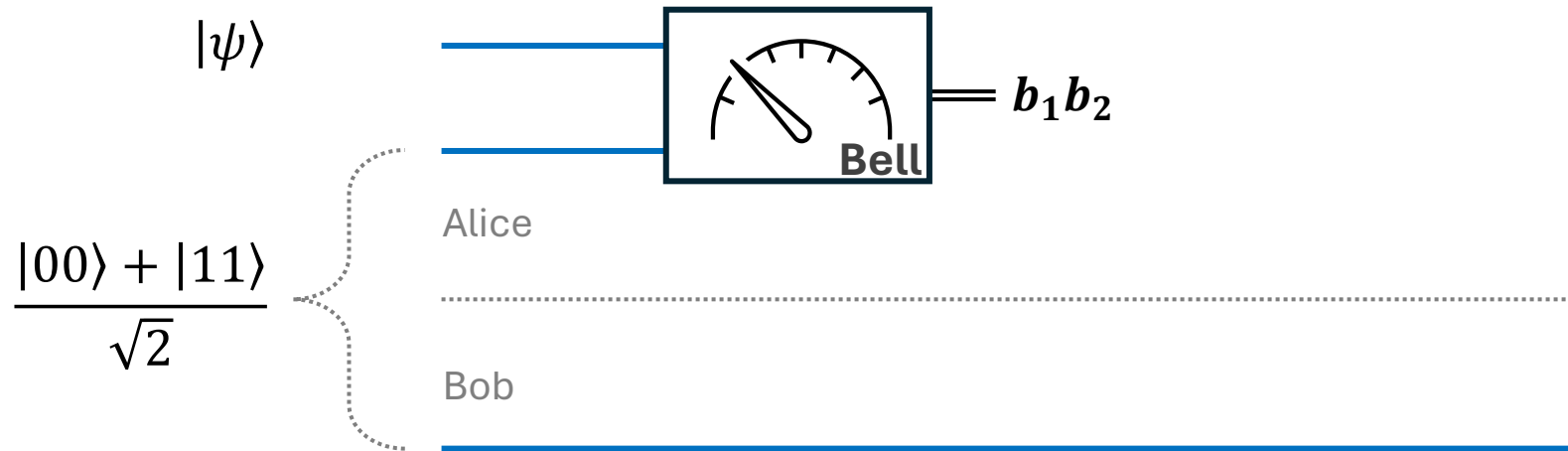
What if Alice does not send b_1 and b_2 to Bob?

Total states:

$$\begin{aligned}
 |\phi_0\rangle &= \sum_{b_1, b_2 \in \{0,1\}} |\beta_{b_1 b_2}\rangle U_{b_2 b_1}^\dagger |\psi\rangle & |\phi_1\rangle &= |\beta_{b_1 b_2}\rangle U_{b_2 b_1}^\dagger |\psi\rangle & |\phi_2\rangle &= |\beta_{b_1 b_2}\rangle \otimes |\psi\rangle
 \end{aligned}$$

Another look at Quantum Teleportation

- Quantum teleportation: Transmit a single-qubit state $|\psi\rangle$.

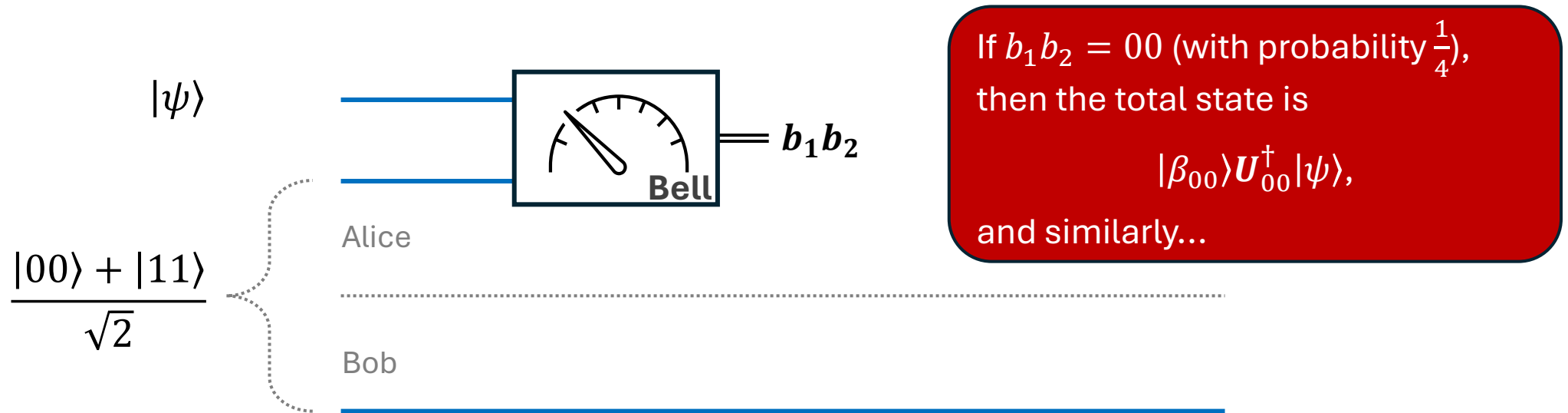


Total states:

$$|\phi_0\rangle = \sum_{b_1, b_2 \in \{0,1\}} |\beta_{b_1 b_2}\rangle U_{b_2 b_1}^\dagger |\psi\rangle \quad |\phi_1\rangle = ?$$

Another look at Quantum Teleportation

- Quantum teleportation: Transmit a single-qubit state $|\psi\rangle$.



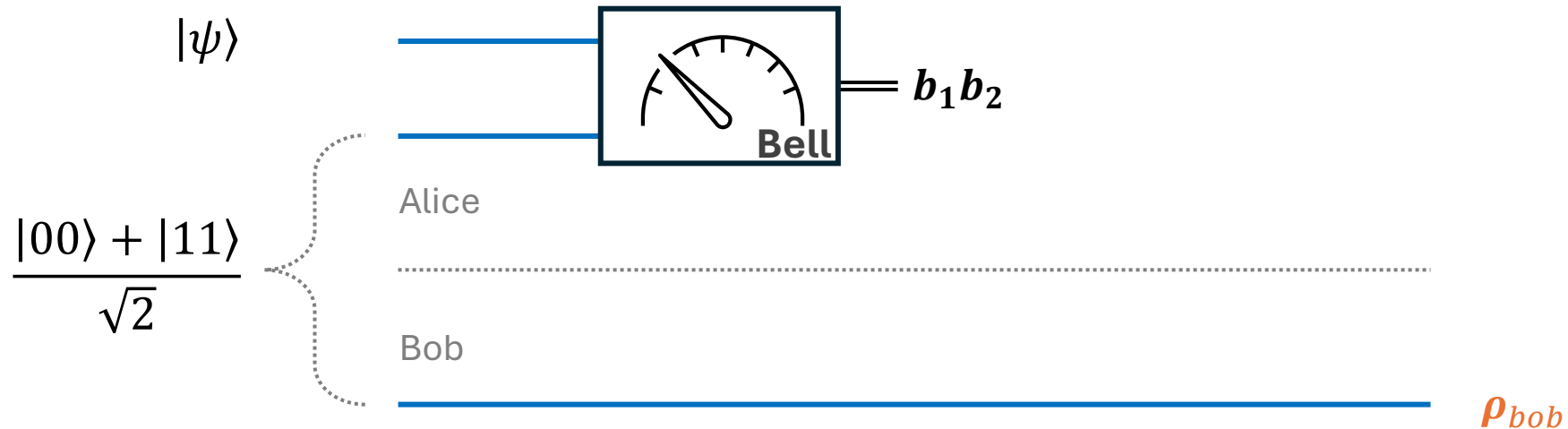
Total states: $|\phi_0\rangle$
 $= \sum_{b_1, b_2 \in \{0,1\}} |\beta_{b_1 b_2}\rangle U_{b_2 b_1}^\dagger |\psi\rangle$

ρ_1 (the density operator on Bob's view after Alice's measurement)

$$= \frac{1}{4} \sum_{b_1, b_2 \in \{0,1\}} |\beta_{b_1 b_2}\rangle U_{b_2 b_1}^\dagger |\psi\rangle \langle \beta_{b_1 b_2}| U_{b_2 b_1} \langle \psi|$$

Another look at Quantum Teleportation

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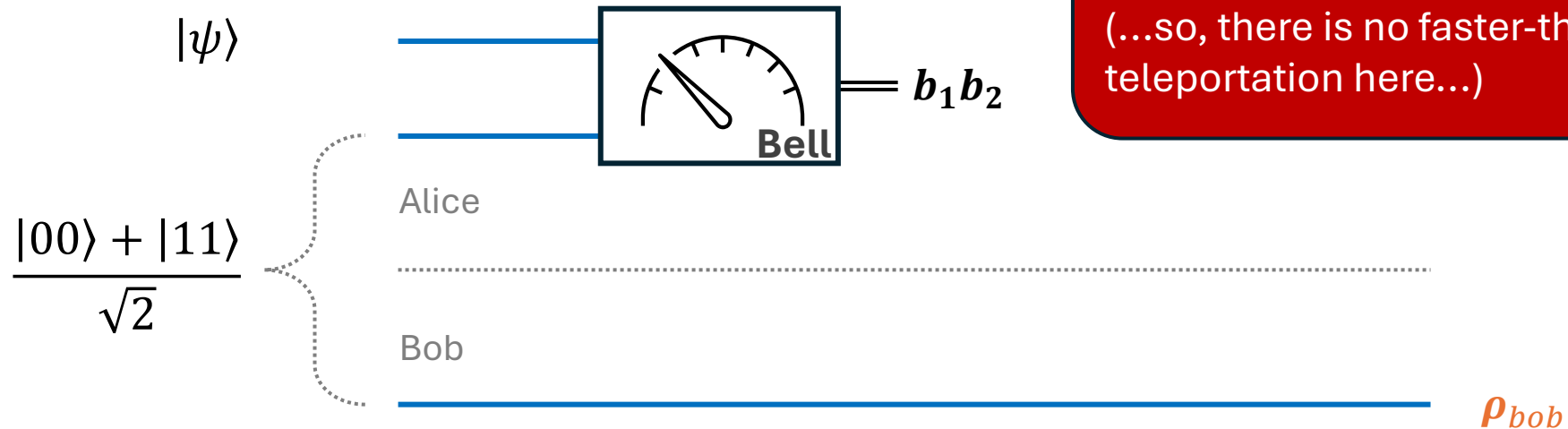


$$\rho_{bob} = \text{tr}_{alice}(\rho_1) = \text{tr}_{alice} \left(\frac{1}{4} \sum_{b_1, b_2 \in \{0,1\}} |\beta_{b_1b_2}\rangle U_{b_2b_1}^\dagger |\psi\rangle \langle \beta_{b_1b_2}| U_{b_2b_1} |\psi\rangle \right)$$

$$= I/2$$

Another look at Quantum Teleportation

- Quantum teleportation: Transmit a single-qubit state $|\psi\rangle$.



If Alice does not send b_1b_2 to Bob, then Bob learns **nothing** about $|\psi\rangle$ (...so, there is no faster-than-light teleportation here...)

$$\rho_{bob} = \text{tr}_{alice}(\rho_1) = \text{tr}_{alice} \left(\frac{1}{4} \sum_{b_1, b_2 \in \{0,1\}} |\beta_{b_1b_2}\rangle U_{b_2b_1}^\dagger |\psi\rangle \langle \beta_{b_1b_2}| U_{b_2b_1} |\psi\rangle \right)$$

$$= I/2 \text{ (independent of } |\psi\rangle \text{)!}$$

Reference

- **[NC00]:** Sections 2.4, 8.3.1, and 9.2.1
- **[KLM07]:** Section 3.5
- **Purification:** [NC00, Section 2.5] and [KLM07, Section 3.5.2]

Next Topics

- Quantum key distribution
- Quantum money
- Summary of this course

Trace Distance

- Let D_1 and D_2 be two probability distributions defined over the same probability space
- Define $D(m) = p_m$, where p_m is the probability that the sample is m
- Trace distance between D_1 and D_2 :

$$\mathbf{TD}(D_1, D_2) := \frac{1}{2} \sum_m |D_1(m) - D_2(m)|$$

- Trace distance **measures how close the two distributions are.**
- Example: Biased coin vs Fair coin

Trace Distance

- Let ρ_1 and ρ_2 be two density operators
- Trace distance between ρ_1 and ρ_2 :

$$\mathbf{TD}(\rho_1, \rho_2) := \frac{1}{2} \mathbf{tr}(|\rho_1 - \rho_2|)$$

- $\mathbf{TD}(\rho_1, \rho_2)$ bound the trace distance between the measurement distributions of ρ_1 and ρ_2 .
 - Namely, let D_1 and D_2 be the measurement distributions of ρ_1 and ρ_2 , respectively.
 - Then $\mathbf{TD}(D_1, D_2) \leq \mathbf{TD}(\rho_1, \rho_2)$ (regardless of what measurement basis we choose...)